Development of a human eye model incorporated with intraocular scattering for visual performance assessment

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Abstract. A biometry-based human eye model was developed by using the empirical anatomic and optical data of ocular parameters. The gradient refractive index of the crystalline lens was modeled by concentric conicoid isoindical surfaces and was adaptive to accommodation and age. The chromatic dispersion of ocular media was described by Cauchy equations. The intraocular scattering model was composed of volumetric Mie scattering in the cornea and the crystalline lens, and a diffuse-surface model at the retina fundus. The retina was regarded as a Lambertian surface and was assigned its corresponding reflectance at each wavelength. The optical performance of the eye model was evaluated in CodeV and ASAP and presented by the modulation transfer functions at single and multiple wavelengths. The chromatic optical powers obtained from this model resembled that of the average physiological eyes. The scattering property was assessed by means of glare veiling luminance and compared with the CIE general disability glare equation. By replacing the transparent lens with a cataractous lens, the disability glare curve of cataracts was generated to compare with the normal disability glare curve. This model has high potential for investigating visual performance in ordinary lighting and display conditions and under the influence of glare.

Keywords: human eye model; gradient index; chromatic dispersion; scattering model; disability glare; veiling luminance.

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1 Introduction

The anatomic and optical properties of human eyes have long been important issues in the fields of ophthalmology, optics, and vision science. Accurate modeling of the human eyes provides a powerful tool to predict visual performance and facilitate ophthalmic surgeries. Dozens of schematic eye models mimicking the optical structure of the eye have been proposed since the beginning of the 20th century. Early models by Helmholtz, Gullstrand, Emsley, and Le Grand have been widely used for first-order calculations. Due to limited ray-tracing computational power and inadequate biological measurements at the time, these models simplified the optical system of human eyes by reducing the number of ocular surfaces, approximating ocular surfaces with spherical surfaces, or using a homogeneous refractive index for the crystalline lens. As a result, they exhibited departures from experimental data in describing on-axis aberration properties. To increase the consistency with experimental data, aspheric surfaces and a lens with a non-homogeneous refractive index have been incorporated in subsequent models. Asphericities were first introduced to ocular surfaces by Lotmar into the anterior surface of the cornea and the posterior surface of the crystalline lens. This model was able to predict the spherical aberration with a homogeneous crystalline lens but required a hypothetic lens with seven shells of stepped refractive indices to describe off-axis aberrations. Navarro et al. proposed an accommodation-dependent model with four centered aspheric surfaces, including anterior and posterior surfaces of the cornea and crystalline lens, and were able to reproduce the on-axis optical performance. They also computed the wavelength-dependent refractive indices by using the Herzberger formula to fit the experimental data in chromatic aberration. However, they used a homogeneous crystalline lens and thus limited the applicability in predicting off-axis aberrations.

Two approaches have been taken in accounting for the non-homogeneity of the crystalline lens, namely a shell model that has a finite number of concentric shells with stepped refractive indices and a gradient-index (GRIN) model that has a smooth index variation usually described by a set of equations. It has been reported that the non-continuous structure in the crystalline lens generates multiple focal planes and hence a GRIN lens model shows better promise. An adaptive GRIN lens model was first introduced by Blaker since then several models have been proposed to match the biological data measured by ever-advancing technologies. In particular, Liou and Brennan performed an extensive search on biometric data to select the ocular parameters in their schematic eye. They incorporated the nasal pupil shift and the visual-axis departure from the optical axis to better resemble the optical realities. They also proposed a GRIN lens model in a parabolic distribution with respect to the axial and radial coordinates. Goncharov and Dainty combined Navarro’s model for off-axis aberrations and Thibos’ on-axis chromatic model and a GRIN lens model of 4th-order polynomials to construct a wide-field age-dependent schematic eye model. Navarro et al. formulated an adaptive GRIN lens model with concentric conicoid isodinical surfaces to include age and accommodation factors. These GRIN models can be integrated into existing schematic eyes for visual simulations.

Another important factor that affects visual quality is light scattering in the eyes. Light scatter in the ocular structure causes
image contrast reductions on fovea, especially when off-axis glare sources are present.\cite{15} Scattering occurs at the iris, retina, and ocular interfaces, in the ocular media, and from cataract and normal aging of the crystalline lens. The major contribution owes to the large protein molecules inside the intraocular media. The sources of scatter contributions in the human eye are as follows: the cornea accounts for about 30%, the lens 40%, and the retina approximately 20% in normal individuals.\cite{15} Also, Van den Berg et al.\cite{15} had shown that the iris and sclera also scatter light, with light-colored eyes scattering more than dark-brown eyes. In the crystalline lens, the scattering particle sizes of protein molecules are 1 to 4 microns in diameter,\cite{15} which is comparable to the wavelength of visible light. Mie scattering, which has stronger forward scattering and is approximately wavelength-independent, is the dominating process for particles of this size. In virtue of our experiences in solid-state lighting,\cite{15,15} we are particularly interested in visual performance under various lighting and display conditions. The scattering properties of ocular media are essential for simulating glare effects on visual contrast sensitivity. Intraocular scattering was first introduced by Navarro\cite{15} into schematic eyes but is missing in modern eye models. In order to incorporate intraocular scattering as well as accommodation and age effects, a biometry-based schematic eye model was developed in this research.

2 Principle

Anatomic and optical data of ocular parameters for the averaged relaxed eye, including radii and asphericities of cornea surfaces, thicknesses between interfaces, refractive indices of ocular media, pupil shift towards the nasal, and visual-axis departure from the optical axis were adopted from the biometric data selected by Liou and Brennan.\cite{15} The mathematical GRIN model for crystalline lens was adopted from Navarro et al.\cite{15} in order to simulate the age and accommodation effects. As the shape of the crystalline lens changes with age and accommodation, the thickness and radii of curvatures were adopted from Dubbelman et al.,\cite{15,15} while the surface asphericities were from Navarro et al.\cite{15}

Chromatic dispersion models of homogeneous ocular media were constructed by using the Cauchy equations proposed by Atchison and Smith.\cite{15} The core and cortex refractive indices of the crystalline lens were modified at particular wavelengths according to the same dispersion model and then used to construct the GRIN model.

In order to describe the light scattering behavior in the human eye model, a scattering model in the cornea, crystalline lens, and retina fundus was constructed. The size distribution, volume fraction, and refractive index of the protein molecules inside the normal-aging and cataractous lenses were adopted from Gilliland et al.,\cite{15} and incorporated into the GRIN lens to simulate volumetric scatters. The retina fundus was considered as a Lambertian surface with reflectance values adopted from Zagers et al.\cite{15} Due to the lack of scattering particle data in the cornea, a volumetric scattering model in the cornea was tuned along with the readily-defined scattering models in the crystalline lens and at the retina fundus to match the CIE disability glare formula.\cite{15}

The optical performance of the human eye model was presented by the point spread function (PSF) for an object at infinity and by the modulation transfer function (MTF) at single and multiple wavelengths. The monochromatic MTF was compared with experimental data from Guirao et al.\cite{15} The chromatic optical powers of the real human eye were used to validate the dispersion model. The scattering model was used to simulate the irradiance at the cornea entrance face and the veiling luminance received at the fovea. The ratio of the veiling luminance to the entrance irradiance was then plotted to compare with the CIE disability glare formula. By replacing the transparent lens with a cataractous lens, the disability glare behavior in cataract eyes was predicted.

3 Modeling

3.1 Geometric Optical Parameters

The schematic eye model was first developed in CodeV for evaluating PSFs, MTFs, and dispersion properties without the scattering model. The same model was then constructed in ASAP with the scattering model for evaluating glare luminance. Table 1 shows the geometric optical parameters of the human eye model with an unaccommodated crystalline lens at age 40 and Fig. 1 shows the schematic drawing. The visual axis is tilted with respect to the optical axis towards the nasal side by 5 deg. Surface 3 consists of the pupil aperture and the anterior surface of the crystalline lens. The pupil is displaced nasally by 0.5 mm,\cite{15} while the anterior lens surface is still centered on the optical axis. The surface asphericity is described by the conic constant $Q$ of a conicoid, which is formulated as

$$ (1 + Q)z^2 - 2zR + \omega^2 = 0, $$

where $z$ is the coordinate on the axis of revolution with the origin at the apex, $\omega = \sqrt{x^2 + y^2}$ is the radial distance to the axis, and $R$ is the radius of curvature at the apex. Figure 2 shows the surface types with different values of conic constants. The conic constants for surfaces 1, 2, 3, and 5 in Table 1 are fitted values to

<table>
<thead>
<tr>
<th>Surface</th>
<th>Radius (nm)</th>
<th>Asphericity ($Q$)</th>
<th>Thickness (nm)</th>
<th>Refractive index (at 555 nm)</th>
</tr>
</thead>
<tbody>
<tr>
<td>1</td>
<td>7.77</td>
<td>−0.18</td>
<td>0.50</td>
<td>1.376</td>
</tr>
<tr>
<td>2</td>
<td>6.40</td>
<td>−0.60</td>
<td>3.16</td>
<td>1.336</td>
</tr>
<tr>
<td>3 [pupil]</td>
<td>10.38</td>
<td>−3.88</td>
<td>2.32</td>
<td>$n_{\text{core}}(z, \omega)^a$</td>
</tr>
<tr>
<td>4</td>
<td>−18.40</td>
<td>−1.59</td>
<td>1.55</td>
<td>$n_{\text{cortex}}(z, \omega)^a$</td>
</tr>
<tr>
<td>5</td>
<td>−5.38</td>
<td>−2.30</td>
<td>16.79</td>
<td>1.336</td>
</tr>
</tbody>
</table>

\cite{15}See Eq. 4 and 5 for the GRIN lens model.
the empirical data. Surface 4 is a hypothetic interface of the GRIN lens model as shown in Fig. 3. The radius of curvature surfaces. The GRIN lens model shown in Fig. 3 consists of the anterior and posterior portions are formulated as

\[ n_{\text{ant}}(z, \omega) = n_0 + \delta_n \left[ 1 - \frac{1}{f_{\text{ant}}} \left( \frac{z^2 - 2\Delta_{\text{ant}}}{a_{\text{ant}}^2} + \epsilon_{\text{ant}} \frac{\omega^2}{b_{\text{ant}}^2} \right) \right]^\rho, \]

\[ n_{\text{pos}}(z, \omega) = n_0 + \delta_n \left[ 1 + \frac{1}{f_{\text{pos}}} \left( \frac{z^2 - 2\Delta_{\text{pos}}}{a_{\text{pos}}^2} + \epsilon_{\text{pos}} \frac{\omega^2}{b_{\text{pos}}^2} \right) \right]^\rho, \]

where \( n_0 = 1.418 \) is the highest refractive index at the lens core at 555 nm and \( \delta_n = -0.047 \) is the difference between the core and cortex refractive indices. \( f_{\text{ant}} \) is the thickness of the anterior portion and is three-fifths of the total lens thickness \( f_{\text{ant}} = 0.6f \). Parameters \( a \) and \( b \) are the semi-axes of the conicoid, and \( \epsilon \) controls the shape (+1 for ellipsoids and -1 for hyperboloids). The semi-axes can be calculated from the apex radius \( R \) and the conic constant \( Q \) by \( a = R/(Q + 1) \) and \( b^2 = Ra \). Parameters \( \Delta_{\text{ant}} \) and \( \Delta_{\text{pos}} \) move the coordinate origin from the center of the conicoid to the apex, and satisfy the relations \( \Delta_{\text{ant}} = \epsilon_{\text{ant}}a_{\text{ant}} \) and \( \Delta_{\text{pos}} = 1 - \epsilon_{\text{pos}}a_{\text{pos}} \). Parameters \( f_{\text{ant}} \) and \( f_{\text{pos}} \) are normalization constants so that the values inside the brackets of Eq. (3) are between 0 and 1.

### 3.2 Chromatic Dispersion

The chromatic dispersion properties of ocular media were described by the Cauchy equation with four terms as

\[ n(\lambda) = A + B/\lambda^2 + C/\lambda^4 + D/\lambda^6, \]

where the corresponding coefficients for each ocular medium are listed in Table 2. For the crystalline lens, only the coefficients for the core (high) and cortex (low) refractive indices are provided. The lens high and low values at 555 nm in Atchison and Smith are 1.406 and 1.386, respectively, taken from Gullstrand’s GRIN lens model. Their fitted Cauchy equations are

\[ n_{\text{High}}(\lambda) = 1.389248 + 6.521218 \times 10^3/\lambda^2 - 6.110661 \times 10^6/\lambda^4 + 5.908191 \times 10^{13}/\lambda^6, \]

\[ n_{\text{High}}(555) = 1.406, \]

\[ n_{\text{Low}}(\lambda) = 1.369486 + 6.428455 \times 10^3/\lambda^2 - 6.023738 \times 10^6/\lambda^4 + 5.824149 \times 10^{13}/\lambda^6, \]

\[ n_{\text{Low}}(555) = 1.386. \]

They also found that for a particular ocular medium, the chromatic dispersion in the range of 400 to 900 nm can be obtained by a constant scaling factor for different refractive
indices. That is, if $n(\lambda)_A$ is the refractive index used in a particular eye model at a reference wavelength $\lambda$ and $n(\lambda)_B$ is the refractive index to be used in another eye model, the refractive index at some other wavelength $\lambda'$ is then

$$n(\lambda)_B = n(\lambda)_A \frac{n(\lambda')_B}{n(\lambda')_A}. \quad (5)$$

In other words, we can simply scale the Cauchy Eqs. (1a) and (1b) and use $n_{\text{High}}(555) = 1.418$ and $n_{\text{Low}}(555) = 1.371$ in our model. The coefficients of the scaled Cauchy equations are given in Table 3. At a particular wavelength, the corresponding lens high and low values are used in Eqs. (6a) and (6b) to generate the GRIN lens [i.e., $n = n_{\text{High}}(\lambda)$ and $\delta_n = n_{\text{Low}}(\lambda) - n_{\text{High}}(\lambda)$]. The chromatic optical powers of the eye model are calculated at five wavelengths, 450.0, 486.1, 555.0, 587.6, and 656.3 nm, and compared to the average dispersion properties of physiological human eyes.

### 3.3 Scattering Model in Human Eyes

Intraocular light scattering is an optical phenomenon that may strongly affect the visual performance of human eyes. The scatter contributions are approximately 30% from the cornea, 40% from the lens, and 20% from the retina in normal individuals. The rest are caused by the iris and sclera, with light-colored irises scattering more than dark-brown ones. For simplicity, this study considered a dark iris to eliminate the iris scattering effect. The constructed scattering model of the cornea, lens, and retina fundus are detailed below.

#### 3.3.1 Crystalline lens

Intraocular volumetric scattering due to protein molecules in the crystalline lens was modeled by Mie scattering in ASAP. This requires several modeling parameters, including the particle sizes, the equivalent obscuration area fraction, and the particle refractive index. The multilamellar bodies (MLBs) observed by Gilliland et al. in age-related nuclear cataracts were

### Table 2 Parameter change in the GRIN lens model with age A (in years) and accommodation D (in diopters).

<table>
<thead>
<tr>
<th>Parameter</th>
<th>Equation/Values</th>
</tr>
</thead>
<tbody>
<tr>
<td>Lens thickness</td>
<td>$t = 2.93 + 0.0236A + D(0.058 - 0.0005A)$</td>
</tr>
<tr>
<td>Radius of curvature for the anterior surface</td>
<td>$R_{\text{ant}} = 1/(1/(12.7 - 0.058A) + 0.0077D)$</td>
</tr>
<tr>
<td>Radius of curvature for the posterior surface</td>
<td>$R_{\text{pos}} = -1/(1/(5.9 - 0.013A) + 0.0043D)$</td>
</tr>
<tr>
<td>Asphericity of the anterior surface</td>
<td>$Q_{\text{ant}} = -2.8 + 0.025A - 0.0013A^2 - 0.25D$</td>
</tr>
<tr>
<td>Asphericity of the posterior surface</td>
<td>$Q_{\text{pos}} = -1 - 0.06A$</td>
</tr>
<tr>
<td>Exponent in the GRIN lens mode</td>
<td>$\chi = 1.1 \times 10^{-7} A^4 + 2.85$</td>
</tr>
</tbody>
</table>

### Table 3 Coefficients of the Cauchy equation for the dispersion model of ocular media.

<table>
<thead>
<tr>
<th>Medium</th>
<th>A</th>
<th>B</th>
<th>C</th>
<th>D</th>
<th>$n$ (555 nm)</th>
</tr>
</thead>
<tbody>
<tr>
<td>Cornea</td>
<td>1.361594</td>
<td>6.009687 $\times 10^3$</td>
<td>-6.760760 $\times 10^8$</td>
<td>5.908450 $\times 10^{13}$</td>
<td>1.376</td>
</tr>
<tr>
<td>Aqueous</td>
<td>1.321631</td>
<td>6.070796 $\times 10^3$</td>
<td>-7.062305 $\times 10^8$</td>
<td>6.147861 $\times 10^{13}$</td>
<td>1.336</td>
</tr>
<tr>
<td>Lens</td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>High</td>
<td>1.401204</td>
<td>6.577339 $\times 10^3$</td>
<td>-6.163249 $\times 10^8$</td>
<td>5.959037 $\times 10^{13}$</td>
<td>1.418</td>
</tr>
<tr>
<td>Low</td>
<td>1.354566</td>
<td>6.358419 $\times 10^3$</td>
<td>-5.958111 $\times 10^8$</td>
<td>5.760697 $\times 10^{13}$</td>
<td>1.371</td>
</tr>
<tr>
<td>Vitreous</td>
<td>1.322357</td>
<td>5.560240 $\times 10^3$</td>
<td>-5.817391 $\times 10^8$</td>
<td>5.036810 $\times 10^{13}$</td>
<td>1.336</td>
</tr>
</tbody>
</table>

*Taken from Atchison and Smith.
adopted as the scattering particles in the GRIN lens model. The cataractous lenses in their study were from subjects with nuclear sclerosis cataract grade from 1 to 3 based on a 0 to 4 scale and without diabetes. The sizes of the MLBs had a distribution of $2.13 \pm 0.64 \mu m$ in diameter from microscopic observations of 117 MLBs. They corrected the diameters to $2.7 \pm 0.7 \mu m$ by a factor of $4/\pi$ to account for the probability that the sections were not made through the maximum cross-sections of the particles. The estimated particle densities were $556 \text{ MLBs/mm}^3$ and $4071 \text{ MLBs/mm}^3$ for the normal-aging transparent lenses and the cataractous lenses, respectively. The refractive index of the lipid shell of the MLBs was estimated to be 1.50 by Matsuzaki et al. The thickness of the lipid shell was estimated to range from 15 to 50 nm. A representative refractive index of 1.49 as suggested by Gilliland et al. for the cytoplasm inside the MLBs was used. The refractive index of the GRIN lens model was applied to the surrounding cytoplasm. As the lipid shell and the interior of the MLBs had similar assumed refractive indices, and the shell was as thin as a few tens of nm, the internal refractive index of the MLBs was applied in the simulation. A normal distribution was used for the particle sizes and the diameters of 1.3 and 4.1 $\mu m$, twice the standard deviation from the mean, were assigned to be the $e^{-2}$ points of the normal distribution. The equivalent obscuration area fractions were calculated as $3.18 \times 10^{-3}$ and $2.33 \times 10^{-2}$ for the normal-aging and cataractous lenses, respectively.

### 3.3.2 Retina fundus

Various fundus reflection models proposed since the 1960’s were reviewed in Berendschot et al. The individual anatomical components in ocular fundus exhibited directional reflectance other than perfectly diffuse surfaces. However, when the integral reflectance was considered, *in vitro* measurements of the porcine fundus showed that the fundus behaved generally as a Lambertian reflector within 20% deviations. Also, previous studies had adopted the Lambertian assumption to either report the integral fundus reflectance or to validate the fundus reflection model. Accordingly, the fundus of the human retina was considered as a Lambertian diffusing surface in this study. One ray hitting the retina surface will be reflected in random directions. The reflectance values at visual wavelengths were adopted from Zagers et al. For instance, the reflectance value at 555 nm is about 0.0135.

### 3.3.3 Cornea

Due to the lack of biometric data for the scattering model in the cornea, the scattering in the cornea was also simulated by adding scattering particles. The model has three volumetric scattering factors, including the particle size distribution, the equivalent obscuration area fraction, and the particle refractive index. After the crystalline lens scattering and retina fundus diffusing model was constructed, these parameters were varied until the veiling luminance of the eye model matched the values calculated from the CIE disability glare general formula. The simulation configuration is shown in Fig. 4, where a glare source at infinity was modeled at different angles departed from the visual axis temporally. The entrance pupil of the eye was set at 4 mm in diameter for typical photopic conditions. The veiling luminance received at the fovea $L_{\text{veil}}$ (in cd/m$^2$) was normalized with the irradiance at the cornea entrance face $E_{\text{glare}}$ (in lux) and compared with CIE general disability-glare formula.

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Fig. 4 Illustration for a glare point source at infinity, its irradiance at the cornea, and the veiling luminance of the glare source contributed to the light received at the fovea.

Fig. 5 Schematic of the human eye model in ASAP (left) and the corresponding PSF on the retina for an object at infinity at 555 nm (right). The contours in PSF represent intensity distributions at 75%, 65%, 45%, 35%, 25%, 15%, and 5% of the maximum intensity.
\[ \frac{L_{\text{veil}}}{E_{\text{glare}}} = \frac{10}{\theta^3} + \left( \frac{5}{\theta^2} + 0.1 \frac{p}{\theta} \right) \cdot \left[ 1 + \left( \frac{\text{Age}}{62.5} \right)^4 \right] + 0.0025p, \quad (6) \]

where \( \theta \) is the angle from the visual axis valid in (0.1, 100) degrees, and \( p \) is a parameter for the ocular pigmentation, ranging from \( p = 0 \) for black eyes to \( p = 1.2 \) for very light eyes. A value of \( p = 0 \) was used in this study to simulate dark-brown eyes in Asian population. The flux detected at the fovea was divided by the area of the exit pupil and the solid angle subtended by the fovea with respect to the exit pupil. This transformation provides the veiling luminance in this study.

4 Results

4.1 Monochromatic MTFs

The schematic of the eye model and the corresponding PSF on the retina for an object at infinity with a wavelength of 555 nm are shown in Fig. 5. The light rays travel parallel to the visual axis (5 deg from the optical axis) before entering the eye and are focused on the fovea. Due to the tilting angle of the visual axis and the nasal shift of the pupil, the resulting spot diagram shows apparent rotational asymmetry. Nevertheless, the MTF in red curve shown in Fig. 6 did not exhibit distinct resolutions in the sagittal and tangential directions. The black dotted line shows the diffraction limit, while the dash-dotted and the solid red lines show the sagittal and tangential MTFs of the eye model, respectively. With the scattering model applied in the human eye model, the mean MTF was redrawn in the blue curve. The solid black line shows the fitted curve to the experimental data in vivo in Guirao et al. for the age group of 40 to 50 years old with a 4-mm pupil size at the wavelength of 543 nm. It is observed that the eye model without the intraocular scattering model overestimated the MTF compared to the experimental data from double-pass measurements. The blue curve in Fig. 6 was the MTF with the intraocular scattering model. The results associated with the scattering model will be detailed in Sec. 4.3.
The blue curve in Fig. 6 showed a nice match to the experimental data when the scattering model was applied in the eye model.

### 4.2 Polychromatic MTFs

The chromatic dispersion curves of the ocular media modeled by the Cauchy equations [Eq. (1) and Table 3] are shown in Fig. 6. The curves for the aqueous and vitreous humors almost coincide. For the crystalline lens, the highest and lowest refractive indices required for constructing the GRIN model are plotted. The refractive-index curves of the lens along the optical axis at five wavelengths, 450.0, 486.1, 555.0, 587.6, and 656.3 nm, are shown in Fig. 6. The chromatic optical powers obtained from the eye model are listed in Table 4 and plotted in Fig. 6. The values for the human eye in reality were calculated by the best-fit Cauchy equation to experimental data in Atchison and Smith.

\[
\Delta \Phi(\lambda) = 1.60911 - 6.70941 \times 10^5 / \lambda^2 + 5.59934 \\
\times 10^{11} / \lambda^4 - 5.59998 \times 10^{15} / \lambda^6, \tag{7}
\]

where \( \Delta \Phi(\lambda) \) is the chromatic difference of optical power in Diopters and was set to zero at 590.0 nm. The maximum difference in Table 4 is 0.021 D, which indicates that the model successfully resembles the chromatic dispersion of human eyes. Figure 7 shows the MTF with five equally-weighted wavelengths at 450.0, 486.1, 555.0, 587.6, and 656.3 nm.

### 4.3 Intraocular Scattering and Veiling Luminance

Figure 8 shows the results of the intraocular scattering model. The horizontal axis indicates the angle between the glare source and the visual axis. The effect of entoptic scattering is demonstrated by the veiling luminance at the retina normalized to the glare illuminance at the cornea entrance face. The red, green, and blue curves are plotted from CIE general disability glare equation for 10, 40, and 70 year-old eyes, respectively. The data points with plus marks are the sampling points of the human eye model with a 40-year-old transparent lens. The parameters of the cornea scattering model were varied to have the data points approaching the CIE general disability glare formula. After an extensive parameter search, the cornea scattering model was tuned to have the refractive index of scattering particles about 1.49, the particle size about 7.0 μm, and an obscuration area of 0.0145 mm² per 1 mm³. These values provide an equivalent cornea scattering model in simulating the CIE general disability glare, but do not necessarily correspond to the realistic scattering particles in cornea. The resulted fitting errors were listed in the form of relative errors in Table 5. The scattering model performed reasonably well in reproducing the veiling luminance compared to that calculated from the CIE general disability glare formula.

By replacing the transparent lens with a cataractous lens, five sampling points of \( L_{\text{veil}} / L_{\text{glare}} \) values corresponding to the glare angles from 1 deg to 5 deg were also plotted in Fig. 9. It is observed that the veiling luminance from a cataractous eye is higher than that from a transparent eye. The ratio of the modeled veiling luminance between the cataractous and the transparent eyes was plotted in Fig. 10. The ratio was about 4 to 7, except for the glare angle of 1 deg. When the glare angle was 1 deg, the flux falling at the detecting fovea area would come from two sources, one from the focal spot of the glare source, and the other from intraocular scattering. Hence the ratio of the veiling luminance between the cataractous and the transparent eyes was altered. The study by de Waard et al. showed that intraocular scattering of nuclear cataracts was about 3 to 4 times of intraocular scattering of normal-aging eyes for the glare angle between 5 deg and 20 deg. The ratio obtained from our model is approximately consistent to the numbers obtained by de Waard et al., despite looking at different glare angles. This assumed angular independency of the ratio can be explained by considering the angular inverse-power laws of the veiling luminance of cataractous and transparent eyes.

### Table 4 Chromatic optical powers of the schematic eye at five wavelengths.

<table>
<thead>
<tr>
<th>Wavelength (nm)</th>
<th>450.0</th>
<th>486.1</th>
<th>555.0</th>
<th>587.6</th>
<th>656.3</th>
</tr>
</thead>
<tbody>
<tr>
<td>Chromatic optical-power difference ( \Delta \Phi(\lambda) ) (m⁻¹)</td>
<td>-1.0453</td>
<td>-0.6688</td>
<td>-0.1760</td>
<td>-0.0043</td>
<td>0.2804</td>
</tr>
<tr>
<td>( \Delta \Phi(\lambda) ) calculated from Eq. (7) (m⁻¹)</td>
<td>-1.0243</td>
<td>-0.6602</td>
<td>-0.1754</td>
<td>-0.0043</td>
<td>0.2807</td>
</tr>
<tr>
<td>Absolute error (m⁻¹)</td>
<td>-0.0210</td>
<td>-0.0086</td>
<td>-0.0006</td>
<td>0</td>
<td>-0.0003</td>
</tr>
<tr>
<td>Relative error (%)</td>
<td>2.05</td>
<td>1.30</td>
<td>0.34</td>
<td>0</td>
<td>-0.11</td>
</tr>
</tbody>
</table>

The resolutions in the sagittal and tangential directions have no significant difference but are lower than the monochromatic MTF (Fig. 6) due to chromatic aberrations.

The blue curve in Fig. 6 showed a nice match to the experimental data when the scattering model was applied in the eye model.
The angular dependence of intraocular straylight in cataractous eyes was found to closely resemble the classic Stiles-Holladay approximation ($10/\theta^2$) and hence approximately inversely proportional to the second power of the glare angle. For comparison, the sampling points for the glare angles of 2 deg to 5 deg from our cataractous lens were curve-fitted to an equation of the inverse-power law. The data from the 1-deg glare angle was excluded as the scattered light detected at the fovea was coupled with the focal spot of the glare source. The resulting fitted curve was

$$\frac{L_{veil}}{E_{glare}}_{\text{cataract}} = 44.36 \times \theta^{-1.85}, \tag{8}$$

with an R-squared value, the coefficient of determination in the regression analysis, of about 0.9746. The multiplication constant $\sim 44$ in Eq. (8) was larger than the value 10 in the Stiles-Holladay formula, suggesting a 4-time increase of the veiling luminance induced by the cataracts. The value of the exponent $-1.85$ was comparable to the value $-2$ in the Stiles-Holladay formula, and was smaller than the value $-2.2$ found by Van den Berg for the angular range of 3.5 deg to 28 deg. Nevertheless, the inverse-power characteristic demonstrated a strong angular dependency of intraocular straylight in cataractous eyes. Figure 13 shows the inverse-power fitting result. This formula is specific for the 40-year-old average human

### Table 5

<table>
<thead>
<tr>
<th>$\theta$ (degrees)</th>
<th>1</th>
<th>2</th>
<th>3</th>
<th>4</th>
<th>5</th>
</tr>
</thead>
<tbody>
<tr>
<td>Relative error</td>
<td>-0.136</td>
<td>0.369</td>
<td>0.069</td>
<td>-0.002</td>
<td>-0.089</td>
</tr>
</tbody>
</table>

$\frac{L_{veil}}{E_{glare}}$

Fig. 10 MTF with five equally-weighted wavelengths at 450.0, 486.1, 555.0, 587.6, and 656.3 nm.

Fig. 11 Veiling luminance due to intraocular scattering. With the p value at 0 and three age parameters, the CIE general disability glare formula is plotted in three curves with red for 10, green for 40, and blue for 70 years-old subjects. The data points with plus marks are samples from the 40-year-old human eye model with a transparent lens. The data points with diamond marks are samples from the 40-year-old human eye model with a cataractous lens.

Fig. 12 The ratio of the modeled veiling luminance between the cataractous and the transparent eyes.
eye with a nuclear-sclerosis cataractous lens, clinically graded from 1 to 3 based on a 0 to 4 scale. Formulas for different ages and cataractous opaque percentages could be generated in a similar fashion.

5 Discussion

In our simulation, the scatter contributions from the cornea, lens, and fundus were found approximately independent for the glare angles from 2 deg to 5 deg. That is, we could remove the scattering properties of the lens and fundus, and calculate the flux detected at the fovea to represent the scattering contribution of the cornea. Similar operations can be done for the lens and the fundus. If we added the three individual contributions together, the result was very close to what we could get by having scatters from the three components at the same time.

Table 6 provides the individual contributions from the three components in our scattering model. For the glare angle as small as 1 deg, the contributions from the cornea and the lens are coupled together and hence were not listed. Figure 13 illustrates the composition of the veiling luminance. It can be observed that the flux provided by fundus scattering was relatively small when the glare angle was below 5 deg. Also, flux provided by fundus scattering was approximately constant with respect to the glare angle. The contributions from the cornea and lens both exhibit a decaying characteristic with respect to the glare angle. However, the curve for the lens did not decay as fast as the curve of the CIE 1999 standard. Hence the curve for the cornea model must decay much faster in order to fit the CIE 1999 standard. When we adjusted the scattering particle size and concentration in the cornea, only one combination of these two parameters would provide the appropriate decay feature. The fitted particle size was about 7.0 μm and the obscuration area was 0.0145 mm² per 1 mm³ in the cornea.

The contributions of the three components in our scattering model have shown some discrepancies to the reported 30-40-20% contributions from cornea-lens-retina in the literature data. However, previous studies have been measured at angles larger than 10 deg, while smaller angles below 5 deg were considered in the present study. For Mie scattering from particles larger than the incident wavelength, the scattering at small angles behaves differently from that at angles larger than 10 deg. For the purpose of assessing visual performance under the influence of glare sources, the calculation of the target image and the veiling luminance could be separated based on the assumption of a linear imaging system. The proposed three-component model can be utilized as an equivalent model to simulate scattering behavior at small glare angles, where the experimental data is currently absent. The image formation can be accomplished by the geometric optical eye model, while the scattering model has enabled us to have a modulation transfer function matching the in vivo measurements as in Fig. 6.

The influence of the glare source at large glare angles can be calculated directly by the CIE general disability glare formula and superimposed to the image. This process is much easier for larger glare angles, where the consideration of the realistic

<table>
<thead>
<tr>
<th>Glare angle</th>
<th>$L_{\text{veil}}/E_{\text{glare}}$ from</th>
<th>Total $L_{\text{veil}}/E_{\text{glare}}$</th>
<th>Contribution (%)</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td>Cornea</td>
<td>Lens</td>
<td>Fundus</td>
</tr>
<tr>
<td>2</td>
<td>1.65</td>
<td>1.42</td>
<td>$3.69 \times 10^{-3}$</td>
</tr>
<tr>
<td>3</td>
<td>0.23</td>
<td>0.85</td>
<td>$3.69 \times 10^{-3}$</td>
</tr>
<tr>
<td>4</td>
<td>0.02</td>
<td>0.49</td>
<td>$3.69 \times 10^{-3}$</td>
</tr>
<tr>
<td>5</td>
<td>0.07</td>
<td>0.15</td>
<td>$3.69 \times 10^{-3}$</td>
</tr>
</tbody>
</table>

Fig. 14 Illustration of the veiling luminance composition in linear-scale axes. The sum of the contributions from the cornea, lens and retina fundus equals the CIE general disability glare curve.

Fig. 13 The curve-fitting result of the disability glare in cataractous eyes.

Table 6 Scatter contributions from the cornea, lens, and fundus.
ocular scatterers is not necessary for the calculation of the veiling luminance.

As the proposed model matching the CIE 1999 standard at glare angles below 5 deg, the parameters used for the scattering particles in the cornea did not correspond to the wavelength dependency of cornea scattering reported in previous studies. The in vitro study on rabbit corneas both suggested a strong wavelength dependency of the inverse 4th power, which corresponds to the Rayleigh scattering by particle sizes much smaller than the wavelength of the incident light. On the contrary, the fitted particle diameter of 7.0 μm in our cornea model corresponds to strong forward scattering by the Mie theory. This inconsistency suggested that the combined effect of the cornea and fundus scattering only represent an equivalent model. Although the integral reflectance of the porcine fundus exhibited the Lambertian characteristics and the Lambertian assumption has been adopted in previous studies, the internal scattering from the multilayer fundus to the retina and directly absorbed at the fovea photoreceptors was not considered. Due to the lack of the experimental data on the angular distribution of internal fundus scattering, the Lambertian characteristic was used in the present model. As shown in Fig. 1, the combined effect of the cornea and fundus must provide a strong decaying behavior with respective to the glare angle. By adopting the Lambertian assumption at the fundus, the angular dependency was compensated by the cornea to yield an equivalent model. In order to have a cornea model corresponding to the in vivo Rayleigh scattering characteristic, it would be sensible to hypothesize internal fundus scattering at very shallow angles towards the retina. This will raise the angular dependency as well as the contribution of the fundus scattering to the veiling luminance. Computations by existing fundus reflection models such as those mentioned in Berendschot et al. or direct measurements would be beneficial to provide the necessary angular distribution of the internal fundus scattering.

In addition, in vitro light scattering measurements from donor lenses by van den Berg at angles from 10 deg to 165 deg suggested wavelength dependency of the lens scattering. Two candidate groups of particles were proposed including a small mean particle size of about 20-nm diameter with a 0.023 volume fraction and a larger mean particle size of about 1.4-μm diameter with a 3 x 10^-6 volume fraction. The corresponding scattering would be a mixture of the Rayleigh and Rayleigh-Gans-Debye types and have strong wavelength dependency. As the current study focuses on the disability glare at angles smaller than 5 deg, it is the forward scattering by particles larger than the incident wavelength that play an important role. Hence the contribution of the Rayleigh scattering from particles as small as 20-nm diameter is negligible. For the model to extend into much larger glare angles, Rayleigh scattering by small particles must be taken into account. As discussed earlier, the veiling luminance contributed from larger glare angles can be readily calculated by the CIE 1999 standard. By having a scattering model matching the CIE 1999 standard for the angular range of 1 deg to 5 deg, the current model yields a modulation transfer function matching the in vivo measurements and provides a simulation tool for visual performance assessment.

6 Conclusion

Abundant ocular biometric data, advances in optical simulation software, and accurate prior models in the literature have enabled us to construct a schematic human eye model that resembles physiological eyes with the functions of age and lens accommodation. The visual acuity and chromatic dispersion of the schematic eye model matched the behavior of real human eyes. The effect of the scattering model in an average eye with a transparent crystalline lens was also confirmed by the CIE disability glare general formula. The disability glare curve of cataract eyes was predicted by replacing the transparent lens with a cataractous lens. A cataractous disability glare formula was then generated by curve-fitting. The comparison between transparent eyes and cataract eyes shows that cataract eyes are more sensitive to glare sources, especially at a glare angle of 2 deg to 4 deg. This model is anticipated to become a powerful tool for studying visual performance with various combinations of different objects, light sources and observers with normal or altered conditions.

Acknowledgments

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References